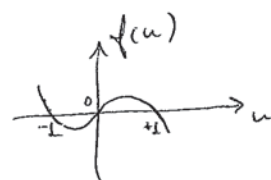


Travelling wave inside a thin channel:

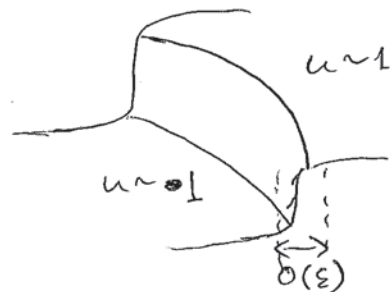
①

Consider the Allen-Cahn model of phase separation:

$$(1) \begin{cases} u_t = \varepsilon^2 \Delta u + f(u), & x \in \Omega \subset \mathbb{R}^2 \\ \partial_n u = 0, & x \in \partial\Omega \\ f(u) = -2u(u-1)(u+1) \end{cases}$$



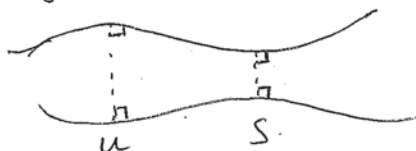
- In the limit $\varepsilon \rightarrow 0$ the solution quickly forms an interface layer of thickness $O(\varepsilon)$, with $u \sim -1$ on one side of the interface and $u \sim 1$ on the other side:



- Once formed, the interface will move slowly (speed $O(\varepsilon^3)$) according to the mean curvature law: the speed in the direction orthogonal to the interface is proportional to the curvature of the interface at that point:



- If Ω is a bounded domain, the interface will move in such a way as to minimize its length.
- The steady state is a straight interface; orthogonal to $\partial\Omega$:



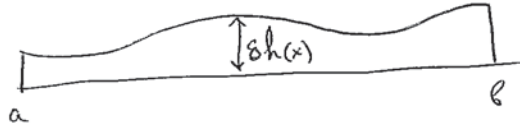
It is stable at the "neck" of the domain, unstable at the "hips".

We consider the case of thin domain:

(2)

$$\Omega = \{ (x, y) : 0 \leq y \leq \delta h(x), a \leq x \leq b \}$$

Where $\delta \ll \varepsilon \ll 1$, $h(x) > 0$.



Then (1) can be reduced to a one-dimensional problem as follows:

Change variables: $z = \frac{y}{\delta h(x)}$, $u(x, y) = U(x, z)$.

Then $u_x = U_x - \frac{y h'}{\delta h^2} U_z = U_x - z \frac{h'}{h} U_z$

$$u_{xx} = U_{xx} - 2z \frac{h'}{h} U_{xz} + \frac{h'^2}{h^2} [z^2 U_{zz} + 2z U_z] - \frac{h''}{h^2} z U_z$$

$$u_y = \frac{1}{\delta h} U_z$$

$$u_{yy} = \frac{1}{\delta^2 h^2} U_{zz}$$

The Boundary conditions are: $\begin{cases} u_x h' \delta - u_y = 0 & \text{at } y = \delta h(x) \\ u_y = 0 & \text{at } y = 0 \end{cases}$

They become: $\begin{cases} (U_x - z \frac{h'}{h} U_z) h' \delta^2 = \frac{U_z}{h}, & \text{at } z = 1 \\ U_z = 0 & \text{at } z = 0 \end{cases}$

Expand $u = u_0(x, z) + \delta^2 u_1(x, z) + \dots$

$O(1)$: $\begin{cases} u_{0,zz} = 0, & z \in (0, 1) \\ u_{0,z} = 0 & \text{at } z = 0, 1 \end{cases} \Rightarrow \boxed{u_0(x, z) = U_0(x)}$

$O(\delta^2)$: $\begin{cases} u_{0,t} = \varepsilon^2 (u_{0,xx} + \frac{1}{h^2} u_{1,zz}) + f(u_0) & (*) \\ u_{1,z} = h h' U_{0,x} & \text{at } z = 1 \\ u_{1,z} = 0 & \text{at } z = 0 \end{cases}$

Integrate (*) w.r.t. z on $z \in [0, 1]$:

$$u_{0,t} = \varepsilon^2 \left(u_{0,xx} + \frac{1}{h^2} \underbrace{\left(u_{1,z} \Big|_0^1 \right)}_{h' h u_{0,x}} \right) + f(u_0)$$

$$\Rightarrow u_{0,t} = \varepsilon^2 \left(u_{0,xx} + \frac{h'(x)}{h(x)} u_{0,x} \right) + f(u_0) \quad (2)$$

Eq'n (2) is the homogenized version of (1) for thin domains, (2) is 1-D pbm.

Drop $u_0 \rightarrow u$, our starting point is :

$$(3) \quad \begin{cases} u_t = \varepsilon^2 \left(u_{xx} + \frac{h'}{h} u_x \right) + f(u), \quad a < x < b \\ u'(a) = 0 = u'(b) \end{cases}$$

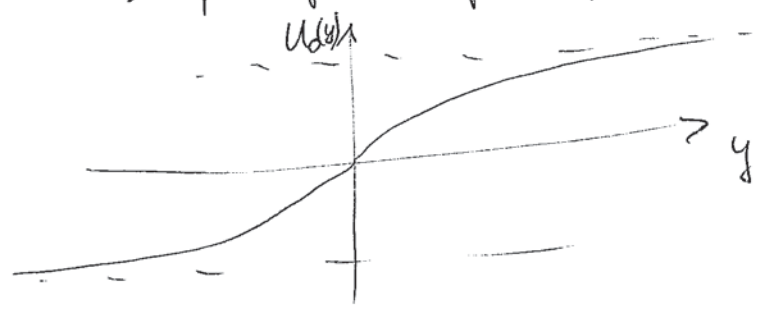
Ansatz : $x = x_0(\varepsilon^2 t) + \varepsilon y$; $u = u_0(y) + \varepsilon u_1 + \dots$

Then $u_x = \frac{1}{\varepsilon} u_y$, $u_{xx} = \frac{1}{\varepsilon^2} u_{yy}$, $\frac{h'(x)}{h(x)} \sim \frac{h'(x_0)}{h(x_0)} + O(\varepsilon)$

$$\Rightarrow O(1) : \begin{aligned} u_t &= -\varepsilon x_0'(\varepsilon^2 t) u_y \\ u_{0,yy} + f(u_0) &= 0 \end{aligned}$$

$$\Rightarrow u_0(y) = \pm \tanh(y)$$

[heteroclinic orbit that describes the shape of interface]



$$O(\epsilon) : \int_{-\infty}^{\infty} U_{0y} (-x_0' U_{0y}) = U_{0y} + U_{0y} f'(U_0) + \frac{h'(x_0)}{h(x_0)} U_{0y}$$

Integrate by parts & use: $(U_{0y})_{yy} + f'(U_0) U_{0y} = 0$:

$$-x_0'(\epsilon t) \int U_{0y}^2 = \frac{h'(x_0)}{h(x_0)} \int U_{0y}^2$$

$$\Rightarrow \boxed{\frac{dx_0}{dt} = -\epsilon^2 \frac{h'(x_0)}{h(x_0)}} \quad (4)$$

- The steady states of ODE (4) correspond to max/min of h ;
- It is stable if $(\frac{h'}{h})' > 0 \Leftrightarrow x_0$ is a min of $h(x)$
- Unstable if x_0 is a max of $h(x)$
- Conclusion: Travelling wave gets "stuck" at the "neck" of the domain (where h has a min).

Numerical example:

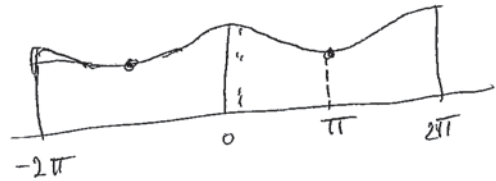
Note that $u(x,t) \sim \pm \tanh\left(\frac{x-x_0(t)}{\varepsilon}\right)$ where

$x_0(t)$ evolves according to (4).

Take $h(x) = 1 + \frac{\cos x}{2}$, $x \in [-2\pi, 2\pi]$

and take i.c.

$$u(x,0) = \tanh\left(\frac{x-0.1}{\varepsilon}\right)$$



Take $\varepsilon = 0.2$; using the software "FlexPDE"

we get the results shown on the next page.

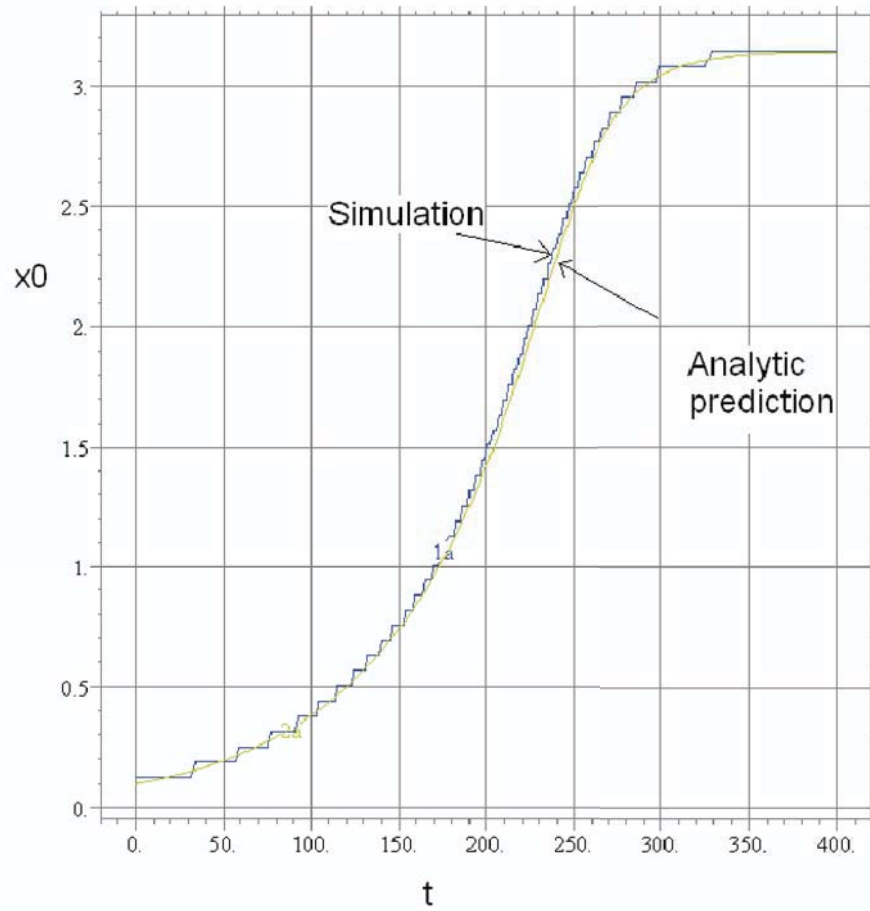
The agreement is very good, as expected,

$$x_0 \rightarrow \pi \text{ as } t \rightarrow \infty$$

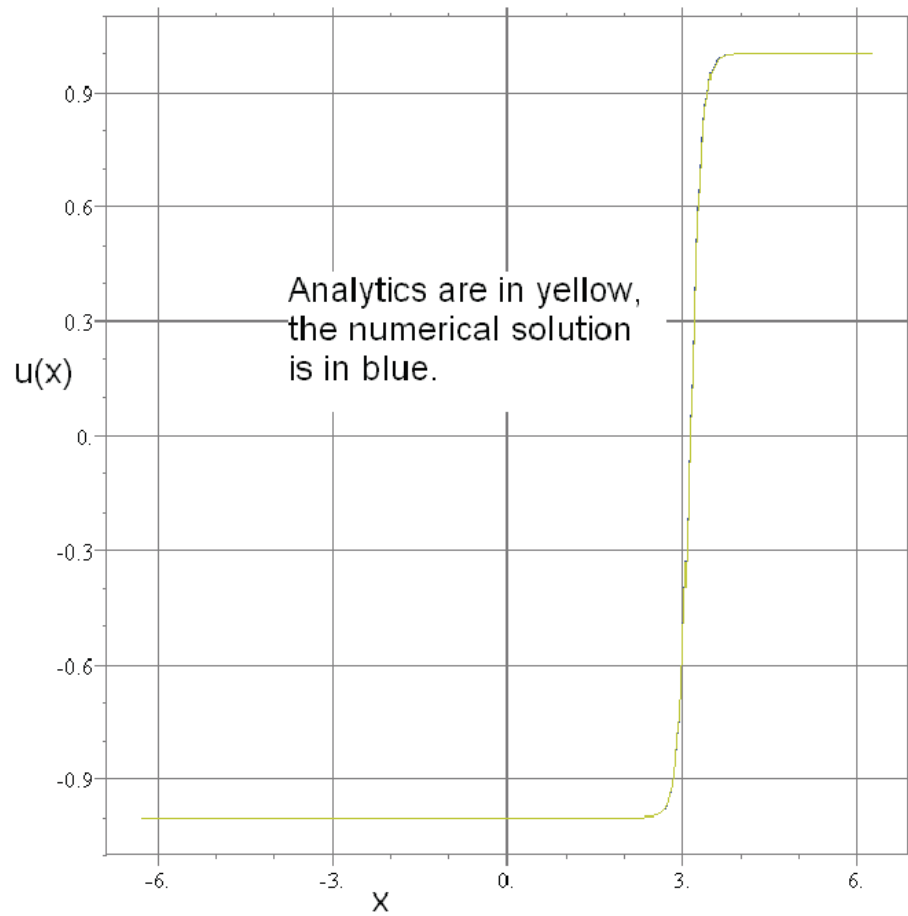
References:

- S. Allen and J.W. Cahn, "A microscopic theory for antiphase boundary motion and its application to antiphase domain coarsening", Acta Metall. 27 (1979), 1084-1095
- J. Rubinstein, P. Sternberg and J.B. Keller, "Fast reaction, slow diffusion, and curve shortening", SIAM J. Appl. Math., Vol. 49 no. 1 (1989), 116-133.

Position of interface vs. time



Interface profile $u(x,t)$ at $t=400$



Further questions

1) Consider the Allen-Cahn system (1):

$$(1) \quad \begin{cases} u_t = \varepsilon^2 \Delta u + f(u), & x \in \Omega, & f(u) = -2u(u^2 - 1) \\ \partial_n u = 0 & \text{on } \partial\Omega \end{cases}$$

where $\Omega \subset \mathbb{R}^3$ is a thin tube, defined as:

$$\Omega = \left\{ (x, y, z) : (y, z) \in \delta D(x), x \in (a, b) \right\}$$

where $\delta \ll \varepsilon \ll 1$ and $D(x) \subset \mathbb{R}^2$ is a 2-d domain that varies smoothly with x .

Show that in the limit $\delta \ll \varepsilon \ll 1$, the system (1)

is approximated by

$$\begin{cases} u_t = \varepsilon^2 \frac{1}{A(x)} (u'(x)A(x))' + f(u), & x \in (a, b) \\ u'(a) = u'(b) = 0 \end{cases}$$

where $A(x) = \text{area}(D(x))$.

2) Consider $u_t = \varepsilon^2 \Delta u + f(u) + \varepsilon g(u), x \in \mathbb{R}^2$

where $f(u) = -2u(u-1)(u+1)$

and $\int_{-1}^1 g(u) = 1$.

a) Find a radially symmetric travelling wave solution of the form:

$$u(x) = U\left(\frac{r - r_0(\varepsilon^2 t)}{\varepsilon}\right), \quad r = |x|.$$

Determine the ODE for r_0 .

b) Does the ODE you found in part (a) have a steady state? Is it stable?