

Bonus problems

(B1)

$$\textcircled{1} \quad \frac{dy}{dx} = \frac{1}{x+y}, \quad y(x_0) = x_0$$

$$\text{let } y = h(x) \Rightarrow \frac{dy}{dx} = h'(x)$$

$$\text{diff. w.r.t } y: \quad 1 = \frac{dx}{dy} h'(x) \frac{dx}{dy}$$

so

$$\frac{dx}{dy} = \frac{1}{h'(x)} = \frac{1}{dy/dx}$$

$$\Rightarrow \frac{dx}{dy} = x+y \quad \leftarrow \text{linear first order eqn}$$

solve by integrating factor ...

$$x(y) = -1 - y + c e^y$$

$$x(y_0) = y(x_0) = y_0 :$$

$$x_0 = -1 - y_0 + c e^{y_0}$$

$$x_0 + 1 + y_0 = c e^{y_0}$$

$$c = e^{-y_0} (1 + x_0 + y_0)$$

so

$$x(y) = -1 - y + e^{y-y_0} (1 + x_0 + y_0)$$

b) vertical tangent in $y(x) \Rightarrow$ slope zero
in $x(y)$

$$\frac{dx}{dy} = -1 + e^{y-y_0} (1 + x_0 + y_0) = 0$$

$$1 + x_0 + y_0 = \underbrace{e^{y_0 - y}}_{\text{always positive}}$$

so we require

$$1 + x_0 + y_0 > 0 \Rightarrow \boxed{y_0 > -1 - x_0}$$

back to $(1+x_0+y_0) = e^{y_0-y^*}$

$$\log[1+x_0+y_0] = y_0 - y^*$$

y_0 $y^* = y_0 - \log(1+x_0+y_0)$

so since $\frac{dx}{dy} = x+y = 0$, $x^* = -y^*$

c) let $x = x^* + \gamma$ $|\gamma| < 1$
 $y = y^* + \epsilon$ $|\epsilon| < 1$

$$x^* + \gamma = -1 - y^* + \epsilon + e^{y^*-y_0+\epsilon} (1+x_0+y_0)$$

since $x^* = -y^*$,

$$\gamma = -1 + \epsilon + e^{y^*-y_0+\epsilon} (1+x_0+y_0)$$

now with $(1+x_0+y_0) = e^{y_0-y^*}$,

we have

$$\gamma = -1 - \frac{3}{2}\epsilon + \epsilon^3$$

has

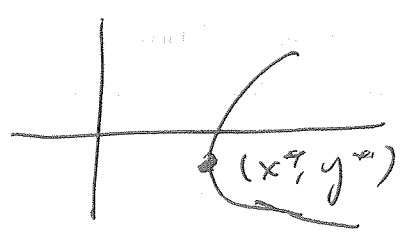
$$\gamma = -1 - \frac{3}{2}\epsilon + 1 + \epsilon + \frac{\epsilon^2}{2}$$

so

$$\gamma = \frac{\epsilon^2}{2} > 0$$

since we have only 1 saddle, and we know it bends toward positive x ($\gamma > 0$), we conclude that y only exists for $x > x^*$.

d) we must have that $x(y)$ looks like



so for $y_0 > y^*$, $y \rightarrow +\infty$ as $x \rightarrow +\infty$.

and for $y_0 < y^*$, $y \rightarrow -\infty$ as $x \rightarrow +\infty$.

for $y_0 > y^*$, and $y \rightarrow +\infty$, we have

that $|e^{y-y_0}| \gg |y|$. So the dominant balance is between x and e^{y-y_0} .

$$\Rightarrow x \sim e^{y-y_0} (1+x_0+y_0)$$

$$\frac{x}{1+x_0+y_0} \sim e^{y-y_0}$$

$$\log \left(\frac{x}{1+x_0+y_0} \right) \sim y - y_0$$

so

$$\boxed{y \sim y_0 + \log \left(\frac{x}{1+x_0+y_0} \right)}$$

as $x \rightarrow +\infty$.

for $y_0 \ll y^*$ with $y \rightarrow \infty$ as $x \rightarrow +\infty$,

$|e^{y-y_0}| \ll |y|$ so dominant balance

is

$x \sim -1 - y$ so $y \sim -1 - x$ as $x \rightarrow \infty$

e) now with $|+x_0 + y_0| < 0$, and $x \rightarrow -\infty$,

if the dominant balance is between

x and y , then ~~$y \sim x$~~ $y \sim -x - 1$

as $x \rightarrow -\infty$. But then $|e^{y-y_0}| \gg |y|$, so

we must have

$x \sim -e^{y-y_0} |1+x_0+y_0|$

$\frac{+|x|}{|1+x_0+y_0|} \sim e^{y-y_0}$

then

$$y \sim y_0 + \log \left| \frac{x}{1+x_0+y_0} \right| \quad \text{as } x \rightarrow \infty$$

as $x \rightarrow +\infty$, we must have instead

$$x \sim -1 - y$$

$$\text{so } y \sim -1 - x \quad \text{as } x \rightarrow +\infty$$

$$\textcircled{2} \text{ a) } u'' = u \quad u = u(x)$$

$$u' u'' = u' u$$

$$\frac{1}{2} \frac{d}{dx} (u')^2 = \frac{1}{2} \frac{d}{dx} u^2$$

$$(u')^2 = u^2 + c$$

$$u' = \sqrt{u^2 + c} \quad \Rightarrow \quad \frac{du}{dx} = \sqrt{u^2 + c}$$

$$\int \frac{du}{\sqrt{u^2+c}} = \int dx$$

to integrate LHS, let $u = \sqrt{c} \sinh \theta$

$$du = \sqrt{c} \cosh \theta d\theta$$

$$\int \frac{du}{\sqrt{u^2+c}} = \int \frac{\sqrt{c} \cosh \theta}{\sqrt{c^2 \sinh^2 \theta + c}} d\theta$$

$$= \int \frac{\sqrt{c} \cosh \theta}{\sqrt{c} \sqrt{\sinh^2 \theta + 1}} d\theta = \int \frac{\sqrt{c} \cosh \theta}{\sqrt{c} \cosh \theta} d\theta$$

$$\Rightarrow \int d\theta = \int dx$$

So $\theta = x + k$ for constant k .

then we have

$$u = \sqrt{c} \sinh(x+k)$$

now we use

$$\sinh(x+y) = \sinh x \cosh y + \sinh y \cosh x$$

so we have

$$u = \sqrt{c} [\sinh x \cosh k + \sinh k \cosh x]$$

$$= \sqrt{c} \cosh k \sinh x + \sqrt{c} \cosh x \sinh k$$

$$\text{let } \sqrt{c} \cosh k = A, \quad \sqrt{c} \sinh k = B$$

then

$$u = A \sinh x + B \cosh x$$

as required.

b) $u'' = -u$

following the same procedure, we have

$$\frac{du}{dx} = \sqrt{c-u^2}$$

$$\int \frac{du}{\sqrt{c-u^2}} = \int dx$$

let $u = \sqrt{c} \sin \theta$

$$du = \sqrt{c} \cos \theta d\theta$$

so

$$\int \frac{du}{\sqrt{c-u^2}} = \int \frac{\sqrt{c} \cos \theta d\theta}{\sqrt{c-c \sin^2 \theta}}$$

$$= \int d\theta$$

$$\Rightarrow \int d\theta = \int dx$$

$$\theta = x + k.$$

then

$$u = \sqrt{c} \sin(\theta + k) = \sqrt{c} [\sin \theta \cos k + \sin k \cos \theta]$$

$$= \sqrt{c} \cos k \sin x + \sqrt{c} \sin k \cos x$$

then with

$$A = \sqrt{c} \cos k, \quad B = \sqrt{c} \sin k,$$

we have

$$u = A \sin x + B \cos x \quad \text{as required.}$$

$$c) \quad u'' + \gamma u' + u = 0$$

$$\text{let } u = e^{-\frac{\gamma}{2}x} v$$

$$u' = -\frac{\gamma}{2} e^{-\frac{\gamma}{2}x} v + v' e^{-\frac{\gamma}{2}x}$$

$$u'' = \frac{\gamma^2}{4} e^{-\frac{\gamma}{2}x} v - \frac{2\gamma}{2} e^{-\frac{\gamma}{2}x} v' + v'' e^{-\frac{\gamma}{2}x}$$

so

$$v'' - \frac{2\gamma}{2} v' + \frac{\gamma^2}{4} v - \frac{\gamma^2}{2} v + \cancel{\gamma v'} + v$$

$$v'' + \left(1 - \frac{\gamma^2}{4}\right) v = 0.$$

$$\text{let } \mu = 1 - \frac{\gamma^2}{4} > 0$$

then

$$v'' + \Gamma v = 0$$

let $x = y/\sqrt{\Gamma}$ and ~~v(x)~~ $v(x) = w(y)$

then

$$\frac{dv}{dx} = \sqrt{\Gamma} \frac{dw}{dy} \Rightarrow \frac{d^2v}{dx^2} = \Gamma \frac{d^2w}{dy^2}$$

then

$$\frac{d^2w}{dy^2} + w = 0 \quad \text{as desired.}$$

d) $\epsilon^2 u'' - mu + Au^2 = 0$; $u \rightarrow 0$ as $|x| \rightarrow \infty$.

$$\frac{\epsilon^2}{m} u'' - u + \frac{A}{m} u^2 = 0$$

let $x = \frac{\sqrt{m}}{\epsilon} y$, $u(x) = v(y)$

then $\frac{d^2}{dx^2} \rightarrow \frac{m}{\epsilon^2} \frac{d^2}{dy^2}$

so

$$v'' - v + \frac{A}{m} v^2 = 0.$$

now let $v = aw$

$$aw'' - aw + a^2 \frac{A}{m} w^2 = 0.$$

choose $a^2 \frac{A}{m} = a \Rightarrow a = \frac{m}{A}$

therefore

$$w'' - w + w^2 = 0, w \rightarrow 0 \text{ as } |y| \rightarrow \infty.$$

e) since $w \rightarrow 0$ as $|y| \rightarrow \infty$, we have

$$w^2 \ll w \text{ as } |y| \rightarrow \infty.$$

$$\Rightarrow w'' \approx w$$

so we have $w \sim ce^{-|y|}$ as $|y| \rightarrow \infty$.
or $u \sim \hat{c} e^{-\frac{\sqrt{m}}{\epsilon}|x|}$ as $|x| \rightarrow \infty$.

f) $w'' - w + w^2 = 0$

$$w'' = -w^2 + w$$

$$w' w'' = w' [w - w^2]$$

$$\frac{1}{2} \frac{d}{dy} (w')^2 = \frac{d}{dy} \left(\frac{w^2}{2} - \frac{w^3}{3} \right)$$

$$\frac{(w')^2}{2} = \frac{w^2}{2} - \frac{w^3}{3} + C$$

So $w, w' \rightarrow 0$ as $|y| \rightarrow \infty$, $C = 0$.

$$(w')^2 = w^2 - \frac{2}{3} w^3 = w^2 \left(1 - \frac{2}{3} w \right)$$

$$w' = w \sqrt{1 - \frac{2}{3} w}$$

$$\int \frac{dw}{w \sqrt{1 - \frac{2}{3} w}} = \int dy$$

$$\text{let } w = \frac{3}{2} \operatorname{sech}^2 \theta$$

$$dw = -3 \operatorname{sech}^2 \theta \tanh \theta d\theta$$

$$\int \frac{dw}{w \sqrt{1 - \frac{2}{3}w}} = \int \frac{-3 \operatorname{sech}^2 \theta \tanh \theta}{\frac{3}{2} \operatorname{sech}^2 \theta \sqrt{1 - \operatorname{sech}^2 \theta}} d\theta$$

now

$$1 - \operatorname{sech}^2 \theta = \tanh^2 \theta$$

$$2 \int \frac{dw}{w \sqrt{1 - \frac{2}{3}w}} = - \int 2 d\theta$$

$$\Rightarrow -2\theta = y + C$$

take $C = 0$ to center maximum at $y = 0$.

then $\theta = -\frac{y}{2}$.

but $\operatorname{sech}^2\left(-\frac{y}{2}\right) = \operatorname{sech}^2\left(\frac{y}{2}\right)$

$$\Rightarrow \boxed{w = \frac{3}{2} \operatorname{sech}^2\left(\frac{y}{2}\right)}$$

g) now $v = \frac{m}{A} w = \frac{3m}{2A} \operatorname{sech}^2\left(\frac{y}{2}\right)$

and

$$x = \frac{\epsilon}{\sqrt{m}} y$$

$$\Rightarrow u(x) = \frac{3m}{2A} \operatorname{sech}^2 \left(\frac{\epsilon}{2\sqrt{m}} x \right)$$

$$u(x) = \frac{3m}{2A} \operatorname{sech}^2 \left(\frac{\sqrt{m}}{2\epsilon} x \right)$$

3)

$$a) \quad x'' + 2\epsilon x' + x = 0, \quad x(0) = 0, \quad x'(0) = 1$$

$$\text{C.E. : } r^2 + 2\epsilon r + 1 = 0$$

$$r = \frac{-2\epsilon \pm \sqrt{4\epsilon^2 - 4}}{2} = -\epsilon \pm \sqrt{\epsilon^2 - 1}$$

with $\epsilon \ll 1$,

$$r = -\epsilon \pm i\sqrt{1 - \epsilon^2}$$

$$x = e^{-\epsilon t} \left[A \cos \omega t + B \sin \omega t \right], \quad \omega = \sqrt{1 - \epsilon^2}$$

$$x(0) = 0 \Rightarrow A = 0$$

$$x'(0) = \epsilon B \Rightarrow B \omega = 1 \Rightarrow B = \frac{1}{\omega}$$

so

$$x(t) = \frac{1}{\sqrt{1-\epsilon^2}} e^{-\epsilon t} \sin \sqrt{1-\epsilon^2} t$$

b) suppose $\epsilon t \ll 1$, then $e^{-\epsilon t} \sim 1 - \epsilon t + O(\epsilon^2)$

$$(1-\epsilon^2)^{-1/2} \sim 1 + \frac{1}{2} \epsilon^2$$

$$(1-\epsilon^2)^{1/2} \sim 1 - \frac{1}{2} \epsilon^2$$

then $\sin \sqrt{1-\epsilon^2} t \sim \sin \left(t - \frac{\epsilon^2}{2} t \right) \sim \sin t + O(\epsilon^2)$

$$\Rightarrow x(t) \sim (1-\epsilon t) \sin t + O(\epsilon^2) \text{ for } \epsilon t \ll 1$$

the expansion $e^{-\epsilon t} \sim 1 - \epsilon t$ required $\epsilon t \ll 1$

so the approximation holds only for

$$0 < t \ll \frac{1}{\epsilon}$$

c) assume $x(t) \sim x_0(t) + \epsilon x_1(t)$

sub into ODE:

$$x_0'' + \epsilon x_1'' + 2\epsilon (x_0' + \epsilon x_1') + x_0 + \epsilon x_1 = 0$$

$$x_0(0) + \epsilon x_1(0) = 0$$

$$x_0'(0) + \epsilon x_1'(0) = 1$$

collect powers of ϵ :

$$O(1): x_0'' + x_0 = 0, \quad x_0(0) = 0, \quad x_0'(0) = 1$$

$$\Rightarrow x_0(t) = \sin t.$$

$$O(\epsilon): x_1'' + 2x_0' + x_1 = 0, \quad x_1(0) = x_1'(0) = 0.$$

$$x_1'' + x_1 = -2x_0' = -2\cos t$$

$$x_1 = A\cos t + B\sin t - t\sin t.$$

$$x_1(0) = 0 \Rightarrow A = 0$$

$$x_1'(0) = B = 0 \Rightarrow B = 0.$$

$$x_1(t) = -t\sin t$$

$$\Rightarrow x(t) \sim \sin t - \epsilon t \sin t + O(\epsilon^2)$$

(in agreement with part b)

notice: ϵt better be "small"
or else the assumption that
 $\epsilon x_1(t) \ll x_0(t)$ would be
violated.
 $\Rightarrow \epsilon t \ll 1$
for validity.

a) $\epsilon = 0.1 : t_{err} \sim 10.52$

$\epsilon = 0.01 : t_{err} \sim 101.8$

$\epsilon = 0.001 : t_{err} \sim 1000.55$

$\epsilon = 0.0001 : t_{err} \sim 10001.24$

clearly, $t_{err} \sim \frac{1}{\epsilon}$, consistent with our assertion that the asymptotic solution is valid only when $t \ll \frac{1}{\epsilon}$.

here, t_{err} is the first time at which the quantity

$$\frac{|x_e(t) - x_a(t)|}{e^{-\epsilon t}}$$

$$\frac{|x_a(t) - x_\epsilon(t)|}{e^{-\epsilon t}}$$

exceeds unity (some measure of the relative error). Here, $x_a(t)$ is the asymptotic solution from part c), $x_e(t)$ is the exact solution from part a).

(4) $y'' + p(t)y' + q(t)y = f(t), \quad y(0) = \alpha, \quad y'(0) = \beta$

let us assume we have two l.i. solutions of the homog. equation y_1, y_2 .

let $y = y_h + y_p$, where

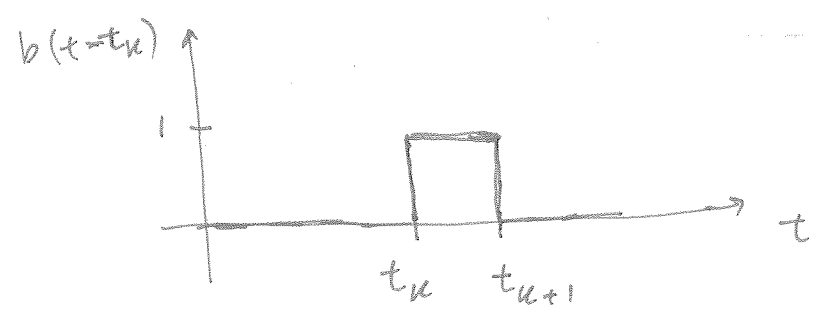
$$\mathcal{L}(y_h) = 0, \quad y_h(0) = \alpha, \quad y_h'(0) = \beta$$

$$\mathcal{L}(y_p) = f(t), \quad y_p(0) = y_p'(0) = 0 \quad (\text{A})$$

(add these together to see that indeed this satisfies original IVP)

write $f(t)$ as

$$f(t) = \lim_{\Delta t \rightarrow 0} \sum_k f(t_k) \frac{b(t-t_k)}{\Delta t} \Delta t; \quad t_k \equiv k\Delta t$$



$$b(t-t_k) = \begin{cases} 1 & t_k < t < t_{k+1} \\ 0 & \text{else} \end{cases}$$

sub into \textcircled{A} :

$$\mathcal{L}(y_p) = \sum_k f(t_k) \frac{b(t-t_k) \Delta t}{\Delta t}; \quad t_k \equiv k \Delta t$$

$$y_p(0) = y_p'(0) = 0$$

by linearity, we may express y_p as

$$y_p = \sum_k y_{pk}(t),$$

where

$$y_{pk}'' + p(t) y_{pk}' + q(t) y_{pk} = f(t_k) b(t-t_k)$$

$$y_{pk}(0) = y_{pk}'(0) = 0$$

notice that $\mathcal{L}(y_{pk}) = 0$, $y_{pk}(0) = y_{pk}'(0) = 0$ for

$0 < t < t_k$. $\Rightarrow y_{pk} = 0$ for $0 \leq t \leq t_k$. A "kick" of

duration Δt and strength $f(t_k) \Delta t$ is applied

at $t = t_k$. To calculate its effect,

integrate ODE from $t = t_k \rightarrow t_{k+1}$ (interval of length Δt):

then we have

$$y_{pk}'(t_{k+1}) - y_{pk}'(t_k) + \int_{t_k}^{t_{k+1}} p(t) y_{pk}' dt + \int_{t_k}^{t_{k+1}} q(t) y_{pk} dt = f(t_k) \Delta t$$

↑
0 since $y_{pk} = 0$ for $t \leq t_k$

notice since $p(t)$ is finite and continuous, we must have

$$y_{pk}'(t_{k+1}) \gg \int_{t_k}^{t_{k+1}} p(t) y_{pk}' dt$$

since the interval of integration is $O(\Delta t)$.

$$\Rightarrow y_{pk}'(t_{k+1}) = O(\Delta t), \quad \int_{t_k}^{t_{k+1}} p(t) y_{pk}'(t) dt = O(\Delta t^2),$$

then since $y_p' = O(\Delta t)$, y_p must be $O(\Delta t^2)$

(since $y_p \approx \Delta t y_p'$) on the interval $t_k < t < t_{k+1}$

⇒ matching orders of Δt , we have

$$y_{pk}'(t_{k+1}) = f(t_k) \Delta t$$

$$y_{pk}(t_{k+1}) = O(\Delta t^2)$$

so we have for y_{pk} :

$$\mathcal{L}(y_{pk}) = 0, \quad y_{pk}(t_{k+1}) = 0$$

$$y_{pk}'(t_{k+1}) = f(t_k) \Delta t$$

then

$$y_{pk}(t) = A_k y_1(t) + B_k y_2(t)$$

$$y_{pk}(t_{k+1}) = 0 \Rightarrow A_k y_1(t_{k+1}) + B_k y_2(t_{k+1}) = 0$$

$$y_{pk}'(t_{k+1}) = f(t_k) \Delta t \Rightarrow A_k y_1'(t_{k+1}) + B_k y_2'(t_{k+1}) = f(t_k) \Delta t$$

$$A_k = \frac{\begin{vmatrix} 0 & y_2(t_{k+1}) \\ f(t_k) \Delta t & y_2'(t_{k+1}) \end{vmatrix}}{\begin{vmatrix} y_1(t_{k+1}) & y_2(t_{k+1}) \\ y_1'(t_{k+1}) & y_2'(t_{k+1}) \end{vmatrix}}, \quad B_k = \frac{\begin{vmatrix} y_1(t_{k+1}) & 0 \\ y_1'(t_{k+1}) & f(t_k) \Delta t \end{vmatrix}}{\begin{vmatrix} y_1(t_{k+1}) & y_2(t_{k+1}) \\ y_1'(t_{k+1}) & y_2'(t_{k+1}) \end{vmatrix}}$$

$$A_k = \frac{-y_2(t_{k+1}) f(t_k) \Delta t}{W_k}$$

$$B_k = \frac{y_1(t_{k+1}) f(t_k) \Delta t}{W_k}$$

$$W_k = \begin{vmatrix} y_1(t_{k+1}) & y_2(t_{k+1}) \\ y_1'(t_{k+1}) & y_2'(t_{k+1}) \end{vmatrix}$$

loosely, we have

$$y = \lim_{\Delta t \rightarrow 0} \sum_k y_k = \lim_{\Delta t \rightarrow 0} \sum_k A_k y_1(t) + B_k y_2(t)$$

$$= y_1(t) \lim_{\Delta t \rightarrow 0} \sum_k A_k + y_2(t) \lim_{\Delta t \rightarrow 0} \sum_k B_k$$

$$= y_1(t) \lim_{\Delta t \rightarrow 0} \sum_k \frac{-y_2(t_{k+1}) f(t_k) \Delta t}{W_k}$$

$$+ y_2(t) \lim_{\Delta t \rightarrow 0} \sum_k \frac{y_1(t_{k+1}) f(t_k) \Delta t}{W_k}$$

in the limit $\Delta t \rightarrow 0$, $t_{k+1} \rightarrow t_k$

then we have

$$y(t) = c_1(t) y_1(t) + c_2(t) y_2(t),$$

$$c_1(t) \equiv - \int \frac{y_2(t) f(t)}{W(y_1, y_2; t)} dt, \quad c_2 \equiv \int \frac{y_1(t) f(t)}{W(y_1, y_2; t)} dt$$
